

Improved Newton Linearization for L^1 -Norm-Type Minimization with **Application to Viscoplastic Fluid Solvers**

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Summary

- ► Our challenge: nonlinear applications modelled by optimization problems with singularities in the Hessian due to a L^1 -norm type term in the objective.
- ► Hessian exhibits a null space upon linearization with Newton's method that is problematic for convergence.
- ► Applications include inverse problems with total variation regularization as well as viscoplastic flows like Earth's mantle convection.
- ▶ We analyze issues with the standard Newton linearization through numerical experiments and theoretically in an abstract setting.
- ▶ We propose an improved linearization based on a perturbation of an otherwise implicitly assumed equality constraint.
- ► We achieve robust and fast Newton convergence independent of the discretization. We further show advantageous theoretical properties of the new linearization.

Introduction in 1 Dimension

Minimization problem Given constants a > 0, $b \in \mathbb{R}$, and $0 \le \theta \le 1$,

find minimizer
$$x^* \in \mathbb{R}$$
 of $\min_{x \in \mathbb{R}} J(x) \coloneqq \frac{a}{1+\theta} |x|_{\epsilon}^{1+\theta} - bx$

Denote $|x|_{\epsilon} := (|x|^2 + \epsilon^2)^{1/2}$. Note that the parameter θ controls the nonlinearity such that Newton's method *converges well* for $\theta = 1$; and Newton's method converges poorly for $\theta \to 0$. The gradient and Hessian of J(x) are

$$g(x) \coloneqq J'(x) = \frac{ax}{|x|_{\epsilon}^{1-\theta}} - b, \quad H(x) \coloneqq J''(x) = \frac{a}{|x|_{\epsilon}^{1-\theta}} \left(1 - (1-\theta) \frac{x^2}{|x|_{\epsilon}^2} \right).$$

Standard Newton linearization One Newton step (simplified, e.g., without backtracking line search) is

solve for
$$\hat{x}$$
 in $H(x) \hat{x} = -g(x)$, update $x \leftarrow x + \hat{x}$,

therefore, we need to solve the linearized equation

$$\frac{a}{|x|_{\epsilon}^{1-\theta}} \left(1 - (1-\theta) \frac{x^2}{|x|_{\epsilon}^2} \right) \hat{x} = -\left(\frac{ax}{|x|_{\epsilon}^{1-\theta}} - b \right) \qquad \text{(standard Newton step)}$$

Perturbed Newton linearization Define for $x, y \in \mathbb{R}$

model error:
$$E(x,y)\coloneqq \frac{x}{|x|_{\epsilon}^{1-\theta}}-y,$$
 model perturbation:
$$D(x,y)\coloneqq x-|x|_{\epsilon}^{1-\theta}y.$$

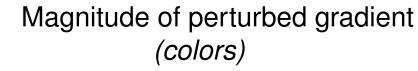
The perturbed Newton linearization for \hat{x} and \hat{y} assumes $D(x,y) \neq 0$, then

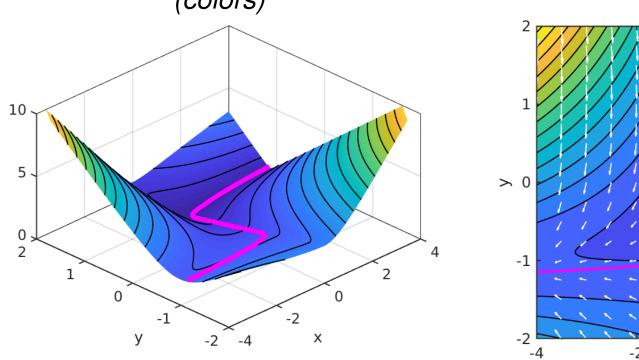
$$a\hat{y} = -(ay - b),$$

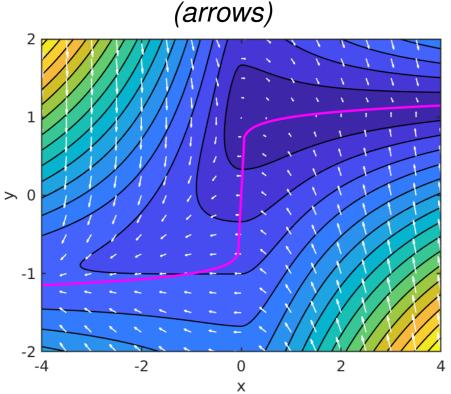
$$\left(1 - (1 - \theta)\frac{xy}{|x|_{\epsilon}^{1+\theta}}\right)\hat{x} - |x|_{\epsilon}^{1-\theta}\hat{y} = -\left(x - |x|_{\epsilon}^{1-\theta}y\right)$$

Substituting for \hat{y} gives the perturbed and reduced Newton linearization for \hat{x}

$$\frac{a}{|x|_{\epsilon}^{1-\theta}} \left(1 - (1-\theta) \frac{xy}{|x|_{\epsilon}^{1+\theta}} \right) \hat{x} = -\left(\frac{ax}{|x|_{\epsilon}^{1-\theta}} - b \right) \qquad \text{(improved Newton step)}$$







Directions of perturbed gradient

The magenta curve in both plots represents the (absolute value of) the unperturbed gradient g(x) in (x,y)-space, meaning E(x,y)=0. Introducing the perturbation, $D(x,y) \neq 0$, relaxes the rapid changes of the unperturbed g(x)around x = 0.

Abstract Derivation of Perturbed Newton Linearizations

Minimization problem Given dimensions $d \in \mathbb{N}$ and $n \in \{1, d, d \times d\}$, domain $\Omega \subseteq \mathbb{R}^d$, and a linear and bounded functional $F: L^2(\Omega)^n \to \mathbb{R}$,

$$\text{find minimizer} \quad U^*:\Omega\to\mathbb{R}^n \quad \text{of} \quad \min_U J(U)\coloneqq \int_\Omega \Phi(U)-F(U)$$

using a generalization of the Huber loss:

$$\Phi(U) \coloneqq \begin{cases} \frac{1}{2} \alpha \, |U|^2 + \beta, & \alpha \, |U| \le \delta, \\ \delta \, |U| - \frac{1}{2} \frac{\delta^2}{\alpha} + \beta, & \text{otherwise} \end{cases}$$

with parameters $\alpha, \delta > 0$ and $\beta \in \mathbb{R}$.

Standard Newton linearization The gradient and Hessian are the 1st- and 2^{nd} -order variations of J:

$$g(U)\tilde{U} \coloneqq \int_{\Omega} \left(\chi \alpha + (1-\chi)\frac{\delta}{|U|}\right) \left\langle U \;, \tilde{U} \right\rangle - F(\tilde{U}), \quad \chi \coloneqq \begin{cases} 1, & \alpha \; |U| \le \delta, \\ 0, & \text{o.w.}, \end{cases}$$

$$\left(H(U)\hat{U} \;, \tilde{U}\right) \coloneqq \int_{\Omega} \left(\chi \alpha + (1-\chi)\frac{\delta}{|U|} \left(\mathbf{I} - \frac{U \otimes U}{|U|^2}\right)\right) \left\langle \hat{U} \;, \tilde{U} \right\rangle.$$

The standard Newton linearization requires to solve for Newton step \hat{U} in

$$\int_{\Omega} \left(\chi \alpha + (1 - \chi) \frac{\delta}{|U|} \left(\mathbf{I} - \frac{U \otimes U}{|U|^2} \right) \right) \left\langle \hat{U}, \tilde{U} \right\rangle = -g(U) \tilde{U}$$

The outer product term is computationally challenging because:

- ► Coefficient $\left(\mathbf{I} \frac{U \otimes U}{|U|^2}\right)$ in Hessian represents an orthogonal projector.
- ▶ Hessian has a zero eigenvalue associated to eigenvector $(1 \chi)U$.

Perturbed Newton linearization Define for $U, S \in L^2(\Omega)^n$

model error:
$$E(U,S) \coloneqq \frac{U}{|U|} - S,$$
 model perturbation: $D(U,S) \coloneqq U - |U| S.$

We augment the previous gradient by a model perturbation, $D(U, S) \neq 0$,

$$g(U,S)\tilde{U} = \int_{\Omega} \left\langle \chi \alpha U + (1-\chi) \delta S \;, \tilde{U} \right\rangle - F(\tilde{U}),$$

$$\left(D(U,S) \;, \tilde{S} \right) = \int_{\Omega} \left\langle U - |U| \; S \;, \tilde{S} \right\rangle,$$

to get the perturbed Newton linearization for step (\hat{U}, \hat{S})

$$\int_{\Omega} \left\langle \chi \alpha \hat{U} + (1 - \chi) \delta \hat{S} , \tilde{U} \right\rangle = -g(U, S) \tilde{U},$$

$$\int_{\Omega} \left\langle \left(\mathbf{I} - \frac{U \otimes S}{|U|} \right) \hat{U} - |U| \hat{S} , \tilde{S} \right\rangle = -\left(D(U, S) , \tilde{S} \right).$$

Here, the dual step has the explicit expression

$$\hat{S} = \frac{U}{|U|} - S + \frac{1}{|U|} \left(\mathbf{I} - \frac{U \otimes S}{|U|} \right) \hat{U}.$$

Therefore, substitution leads to the perturbed and reduced Newton linearization for \hat{U} , where only the outer product term of the Hessian has changed compared to the standard linearization:

$$\int_{\Omega} \left(\chi \alpha + (1 - \chi) \frac{\delta}{|U|} \left(\mathbf{I} - \frac{\boldsymbol{U} \otimes \boldsymbol{S}}{|\boldsymbol{U}|} \right) \right) \left\langle \hat{\boldsymbol{U}} , \tilde{\boldsymbol{U}} \right\rangle = -g(\boldsymbol{U}) \tilde{\boldsymbol{U}}$$

Improvements gained from perturbation

► Perturbation results in model error dependent regularization:

$$\mathbf{I} - \frac{U \otimes S}{|U|} = \mathbf{I} - \frac{U \otimes U}{|U|^2} + \frac{U \otimes E(U, S)}{|U|}.$$

Perturbed linearization acts as a nonlinear preconditioner far from the solution while enabling fast Newton convergence close to solution, since

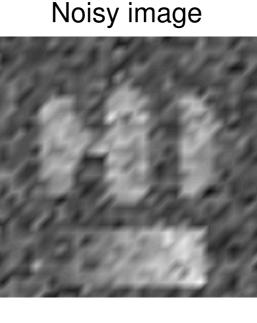
$$\mathbf{I} - \frac{U \otimes S}{|U|} \to \mathbf{I} - \frac{U \otimes U}{|U|^2}$$
 as $E(U, S) \to 0$.

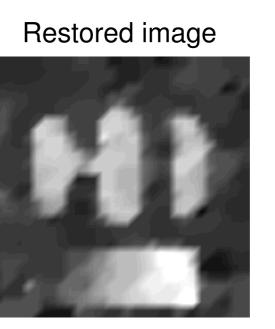
 \blacktriangleright The explicit expression of \hat{S} enables a simple and computationally cheap update of the dual variable S after computing step \hat{U} :

$$S \leftarrow S + \hat{S} = \frac{U}{|U|} + \frac{1}{|U|} \left(\mathbf{I} - \frac{U \otimes S}{|U|} \right) \hat{U}.$$

Image Restoration with Total Variation (TV)

True image





A true image is generated (left), to which Gaussian noise with standard deviation $\sigma=0.2$ is added *(middle)*. The noisy image serves as input data to the image restoration problem with the result shown right.

Minimization problem Given blurry and noisy image data, $d:(0,1)^2 \to \mathbb{R}$, the blurring operator, $B:L^2(\Omega)\to L^2(\Omega)$, $\Omega=(0,1)^2$, and the total variation operator

$$\Phi(\nabla u) = \begin{cases} \frac{1}{2\epsilon} |\nabla u|^2, & |\nabla u| \leq \epsilon, \\ |\nabla u| - \frac{\epsilon}{2}, & \text{otherwise}, \end{cases}$$

with Huber parameters $\alpha = 1/\epsilon$, $\delta = 1$, $\beta = 0$, where $0 < \epsilon \ll 1$; consider

find minimizer
$$u^*:(0,1)^2\to\mathbb{R}$$
 of $\min_u J(u)\coloneqq \frac{1}{2}\left\|Bu-d\right\|_{L^2}^2+\gamma\,\Phi(\nabla u)$

Note: The regularization with $\Phi(\nabla u)$ has the advantage of preserving edges in the restored image u^* . The parameters $\epsilon > 0$ and $\gamma > 0$ regularize the TV operator and the image restoration problem, respectively.

Standard Newton linearization The 1st- and 2nd-order variations of $\Phi(\nabla u)$ are

$$\delta_{u}[\Phi(\nabla u)](\tilde{u}) = \int_{\Omega} \left(\frac{\chi}{\epsilon} + \frac{(1-\chi)}{|\nabla u|}\right) \nabla u \cdot \nabla \tilde{u},$$

$$\delta_{u}\delta_{u}[\Phi(\nabla u)](\tilde{u})(\hat{u}) = \int_{\Omega} \left(\frac{\chi}{\epsilon} + \frac{(1-\chi)}{|\nabla u|}\right) \left(\mathbf{I} - \frac{\nabla u \otimes \nabla u}{|\nabla u|^{2}}\right) \nabla \hat{u} \cdot \nabla \tilde{u}.$$

The standard Newton linearization requires to solve for Newton step \hat{u} in

$$-\nabla \cdot \left[\gamma \left(\frac{\chi}{\epsilon} + \frac{(1-\chi)}{|\nabla u|} \left(\mathbf{I} - \frac{\nabla u \otimes \nabla u}{|\nabla u|^2} \right) \right) \nabla \hat{u} \right] + B^* B \hat{u} = -\gamma \delta_u [\Phi(\nabla u)] - B^* (Bu - d)$$

The left-hand side constitutes a Poisson operator with an anisotropic 2nd-order tensor coefficient.

Perturbed Newton linearization Define for $u \in L^2(\Omega)$, $S \in L^{\infty}(\Omega)$, $||S||_{L^{\infty}} \leq 1$

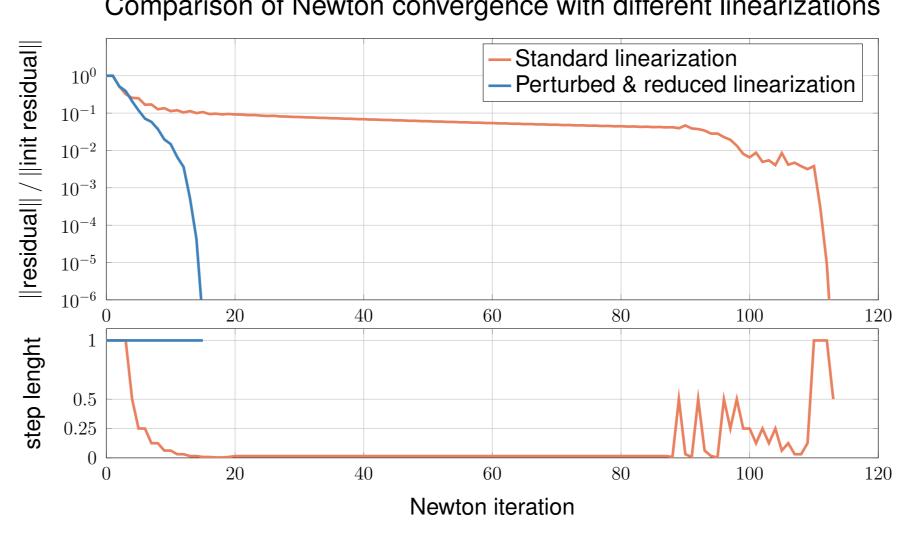
model error: $E(u,S) \coloneqq \frac{\nabla u}{|\nabla u|} - S,$ model perturbation: $D(u, S) := \nabla u - |\nabla u| S$.

The perturbed and reduced Newton linearization for \hat{u} is the system

$$-\nabla \cdot \left[\gamma \left(\frac{\chi}{\epsilon} + \frac{(1 - \chi)}{|\nabla u|} \left(\mathbf{I} - \frac{\nabla u \otimes S}{|\nabla u|} \right) \right) \nabla \hat{u} \right] + B^* B \hat{u} = -\gamma \delta_u [\Phi(\nabla u)] - B^* (Bu - d)$$

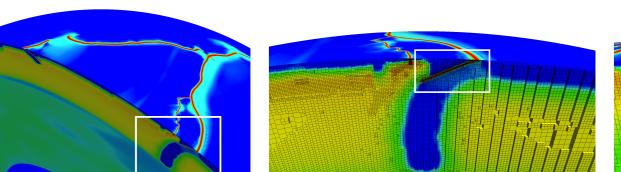
Numerical experiments Comparison of Newton convergence with standard and perturbed linearizations using parameters $\epsilon = 0.01$, $\gamma = 0.002$, and added Gaussian noise to the *true* image with standard deviation $\sigma = 0.2$.

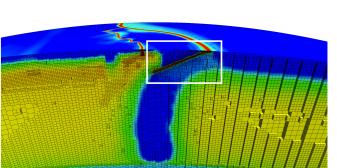
Comparison of Newton convergence with different linearizations

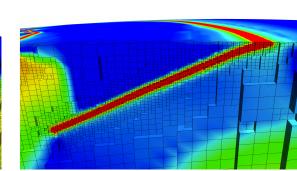


Viscoplastic Stokes Flow

Motivated by rheology of Earth's mantle convection with plate tectonics:







Effective viscosity field (colors) and viscoplastic dynamic weakening of lithosphere in hinge zone (right zoom). (Visualization by L. Alisic)

Minimization problem Given viscosity $\mu > 0$, yield strength $\tau_{\text{yield}} > 0$, the 2nd invariant of the strain rate tensor, $\dot{\varepsilon}_{\text{II}}\coloneqq \frac{1}{\sqrt{2}}\left|\nabla_{\!\!\mathbf{s}} oldsymbol{u}\right|$ with $\nabla_{\!\!\mathbf{s}} oldsymbol{u}\coloneqq \frac{1}{2}(\nabla oldsymbol{u}+\nabla oldsymbol{u}^{\mathsf{T}})$, and

$$\Phi(2\dot{\varepsilon}_{\mathrm{II}}) = \begin{cases} \frac{\mu}{2} (2\dot{\varepsilon}_{\mathrm{II}})^2 + \frac{1}{2} \frac{\tau_{\mathrm{yield}}^2}{\mu}, & \mu(2\dot{\varepsilon}_{\mathrm{II}}) \leq \tau_{\mathrm{yield}}, \\ \tau_{\mathrm{yield}}(2\dot{\varepsilon}_{\mathrm{II}}), & \text{otherwise}, \end{cases}$$

with Huber parameters $\alpha=\mu$, $\delta= au_{\rm yield}$, $\beta=rac{1}{2}rac{\delta^2}{lpha}=rac{1}{2}rac{ au_{
m yield}^2}{\mu}$; consider

$$\text{find minimizer} \ \ \boldsymbol{u}^*:(0,1)^3\to\mathbb{R}^3 \quad \text{of} \quad \min_{\boldsymbol{u}}J(\boldsymbol{u})\coloneqq \int_{\Omega}\Phi(\boldsymbol{u})-\boldsymbol{f}\cdot\boldsymbol{u} \quad \text{s.t.} \quad \nabla\cdot\boldsymbol{u}=0$$

Standard Newton linearization Solve for Newton step (\hat{u}, \hat{p}) in

$$-\nabla \cdot \left[2 \left(\chi \mu + (1 - \chi) \frac{\tau_{\text{yield}}}{\sqrt{2} |\nabla_{\mathbf{s}} \boldsymbol{u}|} \left(\mathbf{I} - \frac{\nabla_{\mathbf{s}} \boldsymbol{u} \otimes \nabla_{\mathbf{s}} \boldsymbol{u}}{|\nabla_{\mathbf{s}} \boldsymbol{u}|^2} \right) \right) \nabla_{\mathbf{s}} \hat{\boldsymbol{u}} \right] + \nabla \hat{p} = -\boldsymbol{r}_{\text{mom}} - \nabla \cdot \hat{\boldsymbol{u}} = -r_{\text{mass}}$$

The residuals of the nonlinear Stokes momentum and mass equations appear on the right-hand side. What plays the role of viscosity in the Newton step is an anisotropic 4th-order tensor.

Perturbed Newton linearization Define for $u \in H^1(\Omega)^d$ and $S \in L^{\infty}(\Omega)^{d \times d}$, $\|\boldsymbol{S}\|_{L^{\infty}} \le 1$

model error: $E(m{u}, m{S}) \coloneqq rac{
abla_{
m s} m{u}}{|
abla_{
m s} m{u}|} - m{S},$

The perturbed and reduced Newton linearization of the momentum equation is

model perturbation: $D(\boldsymbol{u}, \boldsymbol{S}) \coloneqq \nabla_{\!\scriptscriptstyle{\mathrm{S}}} \boldsymbol{u} - |\nabla_{\!\scriptscriptstyle{\mathrm{S}}} \boldsymbol{u}| \boldsymbol{S}$.

$$-\nabla \cdot \left[2\left(\chi \mu + (1-\chi) \frac{\tau_{\text{yield}}}{\sqrt{2} |\nabla_{\!s} \boldsymbol{u}|} \left(\mathbf{I} - \frac{\nabla_{\!s} \boldsymbol{u} \otimes \boldsymbol{S}}{|\nabla_{\!s} \boldsymbol{u}|} \right) \right) \nabla_{\!s} \hat{\boldsymbol{u}} \right] + \nabla \hat{p} = -\boldsymbol{r}_{\text{mom}}$$

Numerical experiments Comparison of Newton convergence with standard and perturbed linearizations using a model problem with a viscosity that incorporates low-viscosity plumes in a high-viscosity background medium

$$\mu(\boldsymbol{x}) \coloneqq (\mu_{\text{max}} - \mu_{\text{min}})(\chi_n(\boldsymbol{x}) - 1) + \mu_{\text{min}}, \quad \boldsymbol{x} \in (0, 1)^3,$$

where $\chi_n(\mathbf{x}) \in [1,2]$ are C^{∞} indicator functions, which accumulate n plumes via products of Gaussians.

Yielding	Mesh	Standard Newton			Perturbed Newton		
volume	level ℓ	It. Newton	#backtr.	It. GMRES	It. Newton	#backtr.	It. GMRES
\sim 45%	4	33	20	1469	10	0	379
\sim 45%	5	36	25	2255	12	0	664
\sim 45%	6	57	49	4255	13	0	876
\sim 65%	4	29	21	1559	18	10	965
$\sim\!\!65\%$	5	37	26	2464	17	9	1245
\sim 65%	6	48	39	3892	20	9	1707
~90%	4	35	25	1505	19	11	872
\sim 90%	5	40	32	2147	21	11	1267
~90%	6	32	21	2312	23	11	1811

References

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- [2] M. Hintermüller and G. Stadler. An infeasible primal-dual algorithm for total variation-based inf-convolutiontype image restoration. SIAM Journal on Scientific Computing, 28(1):1–23, 2006.
- [3] M. Spiegelman, D. A. May, and C. R. Wilson. On the solvability of incompressible Stokes with viscoplastic rheologies in geodynamics. *Geochemistry, Geophysics, Geosystems*, 17(6):2213–2238, 2016.